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Constrained Evolution and Differential Boundary Conditions

**(The sand box for a young
numerical relativist)**

Proposed questions

- How do we define constrained evolution?
- How many boundary conditions are needed for constrained evolution?
- Can there be boundary conditions other than Neumann and Dirichlet?
- How energy estimates motivate boundary conditions?

What is constrained evolution?

$$\begin{aligned}\partial_t^2 u_i &= \partial^j \partial_j u_i \\ \partial^j u_j &= 0\end{aligned}$$

$$\begin{aligned}\partial_t^2 u_i &= \partial^j \partial_j u_i \\ \partial^j u_j &= 0\end{aligned}$$

+ full set of ID
+ full set of BD
=
overdetermined

$$\begin{aligned}\partial_t^2 u_i &= \partial^j \partial_j u_i \\ \partial^j u_j &= 0\end{aligned}$$

+ compatible ID
+ compatible BD
=
hope for constraint preservation

$$\begin{aligned}\partial_t^2 u_i &= \partial^j \partial_j u_i \\ &\quad - \partial_i \partial^j u_j \\ \partial^j u_j &= 0\end{aligned}$$

+ compatible ID
+ ??? many BD
=
well-posed ???

The 1D example

$$\partial_t^2 u = \partial_x^2 u$$

$$\partial_x u = 0$$

How many boundary conditions? (free evolution)

(D.Arnold, N.Tarfulea, A.Alekseenko)

- Consider the constraint quantity $C = \partial_x u$.
- Notice that C satisfies the wave equation $\partial_t^2 C = \partial_x^2 C$.

- Require either

$$C(a) = C(b) = 0$$

$$\partial_x u(a) = \partial_x u(b) = 0$$

or

$$\partial_x C(a) = \partial_x C(b) = 0$$

$$\partial_x^2 u(a) = \partial_x^2 u(b) = 0$$

$$\partial_t^2 u(a) = \partial_t^2 u(b) = 0$$

$$u(a) = u(b) = 0$$

The 1D example

$$\partial_t^2 u = \partial_x^2 u$$

$$\partial_x u = 0$$

How many boundary conditions? (free evolution)

(G. Calabrese, J. Pullin, O. Sarbach, M. Tiglio, O. Reula)

- Consider the constraint quantity $C = \partial_x u$.
- Notice that C satisfies the wave equation $\partial_t^2 C = \partial_x^2 C$.
- Require, for example,

$$\partial_x C(a) = 0, \quad C(b) = 0$$

$$\partial_x^2 u(a) = 0, \quad \partial_x u(b) = 0$$

$$\partial_t^2 u(a) = 0,$$

$$u(a) = u_0(a)t + u_1(a)$$

The 1D example

$$\begin{aligned}\partial_t^2 u &= \partial_x^2 u \\ \partial_x u &= 0\end{aligned}\tag{1}$$

Needs no boundary conditions!

Indeed:

- Replace with

$$\begin{aligned}\partial_t^2 u &= 0 \\ \partial_x u &= 0\end{aligned}\tag{2}$$

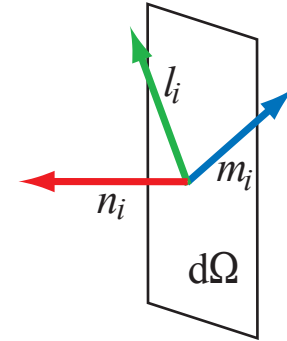
- Integrate (2) from any consistent initial data:

$$u(x, t) = u_0(x)t + u_1(x), \quad \partial_x u_0 = \partial_x u_1 \equiv 0$$

Vector wave equation (constrained evolution)

$$\partial_t^2 u_i = \partial^j \partial_j u_i \quad (3)$$

$$\partial^j u_j = 0$$



How many boundary conditions?

Two, if can eliminate “constraint dependency.”

Subtracting $\partial_i \partial^j u_j$ from (3),

$$\partial_t^2 u_i = 2\partial^j \partial_{[j} u_{i]} \quad (4)$$

$$\partial^j u_j = 0$$

Two boundary conditions:

$$n_{[i} u_{j]} = 0 \quad \text{or} \quad u_i m^i = 0, \quad u_i l^i = 0.$$

Proof of $n_{[i}u_{j]} = 0$: Contract (4) with $\partial_t u^i$, integrate,

$$\int_{\Omega} (\partial_t^2 u_i) \partial_t u^i = \int_{\Omega} 2(\partial^j \partial_{[j} u_{i]}) \partial_t u^i$$

Integrating by parts:

$$\frac{1}{2} \partial_t \left[\|\partial_t u_i\|^2 + 2\|\partial_{[j} u_{i]}\|^2 \right] = \int_{\partial\Omega} (\partial_{[j} u_{i]}) n^{[j} \partial_t u^i]$$

The rest follows from the identity:

$$\|\partial_{(i} u_{j)}\|^2 = \|\partial_{[i} u_{j]}\|^2 + \int_{\partial\Omega} (\partial_j n^i u_i) u^j + \|\partial^i u_i\|^2 - \int_{\partial\Omega} u_i n^i \partial^j u_j$$

(check $u_i m^i = u_i l^i = 0$, thus $(\partial_j n^i u_i) u^j = (\frac{\partial}{\partial n} u_i n^i) u_j n^j = 0$,
since $\partial^j u_j = \frac{\partial}{\partial n} u_i n^i + \frac{\partial}{\partial m} u_i m^i + \frac{\partial}{\partial l} u_i l^i = 0$.)

Inhomogeneous Dirichlet data $n_{[i}u_j] = g_3n_{[i}m_j] + g_2n_{[i}m_j]$
 is equivalent to

$$u_j m^j = g_2, \quad u_j l^j = g_3, \quad \text{implies } \frac{\partial}{\partial n} u_j n^j = -\frac{\partial}{\partial m} g_2 - \frac{\partial}{\partial l} g_3$$

(from $\partial^i u_i \equiv 0$).

Radiation type condition $(\partial_{[i}u_j])n^{[i}\partial_t u^j] \leq 0$, for example,

$$\frac{\partial}{\partial n} u_j m^j + \partial_t u_j m^j = \frac{\partial}{\partial m} u_j n^j, \quad \frac{\partial}{\partial n} u_j l^j + \partial_t u_j l^j = \frac{\partial}{\partial l} u_j n^j.$$

$$\text{(implies)} \quad \partial_t \left(\partial_t u_j n^j + \frac{\partial}{\partial n} u_j n^j \right) = 0$$

(from $\partial^i u_i \equiv 0$, $\frac{\partial}{\partial n} \partial^j u_j = 0$, commuting derivatives and (3))

Differential BCs (conserving $\partial^i u_i = 0$):

$$\partial_t^2 u_i = \partial^j \partial_j u_i$$

give constraint compatible $u_i(0)$, $\partial_t u_i(0)$ and

$$n_{[i} u_{j]} = 0 \quad \text{and} \quad \partial^j u_j = 0$$

Proof. Notice that $C = \partial^j u_j$ satisfies the wave equation $\partial_t^2 C = \partial^j \partial_j C$, with $C(0) = \partial_t C(0) = 0$.

The second boundary condition implies $C|_{\partial\Omega} = 0$. Thus,

$$C \equiv 0, \quad \Rightarrow \quad \partial^j u_j \equiv 0$$

Verify $n_{[i} u_{j]} = 0$ as in the previous example.

More fancy BCs:

$$\partial_t^2 u_i = \partial^j \partial_j u_i$$

give constraint compatible $u_i(0)$, $\partial_t u_i(0)$ and

$$\partial_{[i} u_{j]} n^i = 0 \quad \text{and} \quad \partial^j u_j = 0$$

(Motivated by the energy identity

$$\frac{1}{2} \partial_t \left[\|\partial_t u_i\|^2 + 2 \|\partial_{[j} u_{i]}\|^2 \right] = \int_{\partial\Omega} (\partial_{[j} u_{i]}) n^{[j} \partial_t u^{i]})$$

Reduces to $(\rho = \partial_t u_i(0) n^i t + u_i(0) n^i)$

$$u_i n^i = \rho, \quad \frac{\partial}{\partial n} u_i m^i = \frac{\partial}{\partial m} \rho, \quad \frac{\partial}{\partial n} u_i l^i = \frac{\partial}{\partial l} \rho.$$

Appendix

Linearized BSSN equations

$$\partial_t \varphi = -\frac{1}{6}\kappa + \frac{1}{6}\partial^s \beta_s, \quad \partial_t \alpha = -\kappa, \quad \partial_t \kappa = -\partial^l \partial_l \alpha,$$

$$\partial_t \tilde{\gamma}_{ij} = -2\tilde{A}_{ij} + 2\partial_{(i}\beta_{j)} - \frac{2}{3}\delta_{ij}\partial^s \beta_s,$$

$$\begin{aligned} \partial_t A_{ij} = \frac{1}{2}\partial^l \partial_l \tilde{\gamma}_{ij} + \partial_{(i}\Gamma_{j)} - 2\partial_i \partial_j \varphi - 2\delta_{ij}\partial^l \partial_l \varphi \\ - \partial_i \partial_j \alpha + \frac{1}{3}\delta_{ij}\partial^l \partial_l \alpha, \end{aligned}$$

$$\partial_t \Gamma_i = -\frac{4}{3}\partial_i \kappa + \frac{1}{3}\partial_i \partial^s \beta_s + \partial^l \partial_l \beta_i,$$

Constraint equations:

$$\partial^p \partial^q \tilde{\gamma}_{pq} - 8\partial^l \partial_l \varphi = 0, \quad \text{and/or} \quad \partial^l \Gamma_l - 8\partial^l \partial_l \varphi = 0$$

$$\partial^j A_{ij} - \frac{2}{3}\partial_i \kappa = 0,$$

$$\Gamma_j = \partial^l \tilde{\gamma}_{lj}.$$

Reduction to second order in time

$$\begin{aligned}\partial_t^2 A_{ij} &= \partial^l \partial_l A_{ij} \\ \partial_t^2 k &= \partial^l \partial_l k.\end{aligned}$$

Introduce

$$M_i = \partial^j A_{ij} - \frac{2}{3} \partial_i k$$

Evolution of constraint M_i :

$$\partial_t^2 M_j = \partial^l \partial_l M_j$$

Initial data $M_j(0) = \partial_t M_i(0) = 0$ for compatible data. Thus, $M_i \equiv 0$ as long as

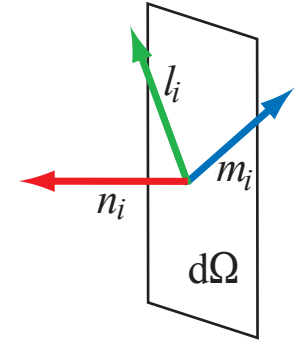
$$\left(\frac{\partial}{\partial n} M_j\right) \partial_t M^j \leq 0 \quad \text{on} \quad \partial\Omega$$

$(\frac{\partial}{\partial n} M_j) \partial_t M^j = 0$ **via the main variables.**

Introduce, orthonormal basis: $n_i, m_i, l_i,$

Rewrite, $\partial^j u_j = \frac{\partial}{\partial n} u_i n^i + \frac{\partial}{\partial m} u_i m^i + \frac{\partial}{\partial l} u_i l^i,$

$$\partial_i \kappa = \frac{\partial}{\partial n} \kappa n_i + \frac{\partial}{\partial m} \kappa m_i + \frac{\partial}{\partial l} \kappa l_i$$



Decompose

$$A_{ij} = A1 (n_{(i} m_{j)}) + A2 (n_{(i} l_{j)}) + A3 (l_{(i} m_{j)}) \\ + A4 (l_i l_j - m_i m_j) + A5 (2n_i n_j - l_i l_j - m_i m_j).$$

Substitute into $M_i = \partial^j A_{ij} - \frac{2}{3} \partial_i k$ and the result into $(\frac{\partial}{\partial n} M_j) \partial_t M^j = 0$

Simplify ...

$$\left(\frac{\partial}{\partial n}M_i\right)(\partial_tM^i)$$

$$= \frac{\partial}{\partial n}\left[\frac{1}{2}\frac{\partial}{\partial m}A1 + \frac{1}{2}\frac{\partial}{\partial l}A2 + 2\frac{\partial}{\partial n}A5 - \frac{2}{3}\frac{\partial}{\partial n}\kappa\right]$$

$$\times \partial_t\left[\frac{1}{2}\frac{\partial}{\partial m}A1 + \frac{1}{2}\frac{\partial}{\partial l}A2 + 2\frac{\partial}{\partial n}A5 - \frac{2}{3}\frac{\partial}{\partial n}\kappa\right]$$

$$+ \frac{\partial}{\partial n}\left[\frac{1}{2}\frac{\partial}{\partial n}A1 + \frac{1}{2}\frac{\partial}{\partial l}A3 - \frac{\partial}{\partial m}A4 - \frac{\partial}{\partial m}A5 - \frac{2}{3}\frac{\partial}{\partial m}\kappa\right]$$

$$\times \partial_t\left[\frac{1}{2}\frac{\partial}{\partial n}A1 + \frac{1}{2}\frac{\partial}{\partial l}A3 - \frac{\partial}{\partial m}A4 - \frac{\partial}{\partial m}A5 - \frac{2}{3}\frac{\partial}{\partial m}\kappa\right]$$

$$+ \frac{\partial}{\partial n}\left[\frac{1}{2}\frac{\partial}{\partial n}A2 + \frac{1}{2}\frac{\partial}{\partial m}A3 + \frac{\partial}{\partial l}A4 - \frac{\partial}{\partial l}A5 - \frac{2}{3}\frac{\partial}{\partial l}\kappa\right]$$

$$\times \partial_t\left[\frac{1}{2}\frac{\partial}{\partial n}A2 + \frac{1}{2}\frac{\partial}{\partial m}A3 + \frac{\partial}{\partial l}A4 - \frac{\partial}{\partial l}A5 - \frac{2}{3}\frac{\partial}{\partial l}\kappa\right](= 0 ??)$$

By a direct observation, either of the two sets of boundary conditions are constraint-preserving:

$$A1 = 0, \quad A2 = 0, \quad \frac{\partial}{\partial n} A3 = 0, \quad \frac{\partial}{\partial n} A4 = 0, \\ \frac{\partial}{\partial n} A5 = 0, \quad \frac{\partial}{\partial n} \kappa = 0,$$

$$\frac{\partial}{\partial n} A1 = 0, \quad \frac{\partial}{\partial n} A2 = 0, \quad A3 = 0, \quad A4 = 0, \\ A5 = 0, \quad \kappa = 0.$$

(the first set eliminates the second multiplier in the first term of $((\partial/\partial n)M^i)(\partial_t M_i) = 0$ and the first multipliers in the second and third terms (by commuting partial derivatives and using evolution eqn. The second set is verified in a similar way.)

Too restrictive!

Differential boundary conditions

Require $M_i = 0$ on $\partial\Omega$:

$$\frac{1}{2} \frac{\partial}{\partial m} A1 + \frac{1}{2} \frac{\partial}{\partial l} A2 + 2 \frac{\partial}{\partial n} A5 - \frac{2}{3} \frac{\partial}{\partial n} \kappa = M_i n^i = 0$$

$$\frac{1}{2} \frac{\partial}{\partial n} A1 + \frac{1}{2} \frac{\partial}{\partial l} A3 - \frac{\partial}{\partial m} A4 - \frac{\partial}{\partial m} A5 - \frac{2}{3} \frac{\partial}{\partial m} \kappa = M_i m^i = 0$$

$$\frac{1}{2} \frac{\partial}{\partial n} A2 + \frac{1}{2} \frac{\partial}{\partial m} A3 + \frac{\partial}{\partial l} A4 - \frac{\partial}{\partial l} A5 - \frac{2}{3} \frac{\partial}{\partial l} \kappa = M_i l^i = 0$$

For example, prescribe $A3$, $A4$, κ , use $M_i = 0$ as the boundary conditions for the rest.

Well-posedness ???

Evolving boundary conditions

Require $\frac{\partial}{\partial n} M_i n^i = 0$, $M_i l^i = 0$, $M_i m^i = 0$. Solve for

$$\begin{aligned} 2 \frac{\partial^2}{\partial t^2} A5 - \left(\frac{\partial^2}{\partial l^2} + \frac{\partial^2}{\partial m^2} \right) A5 - \frac{2}{3} \frac{\partial^2}{\partial t^2} \kappa + \frac{4}{3} \left(\frac{\partial^2}{\partial l^2} + \frac{\partial^2}{\partial m^2} \right) \kappa \\ = \left(\frac{\partial^2}{\partial l^2} - \frac{\partial^2}{\partial m^2} \right) A4 + \frac{\partial}{\partial l} \frac{\partial}{\partial m} A3, \end{aligned}$$

$$\frac{1}{2} \frac{\partial}{\partial n} A1 + \frac{1}{2} \frac{\partial}{\partial l} A3 - \frac{\partial}{\partial m} A4 - \frac{\partial}{\partial m} A5 - \frac{2}{3} \frac{\partial}{\partial m} \kappa = 0$$

$$\frac{1}{2} \frac{\partial}{\partial n} A2 + \frac{1}{2} \frac{\partial}{\partial m} A3 + \frac{\partial}{\partial l} A4 - \frac{\partial}{\partial l} A5 - \frac{2}{3} \frac{\partial}{\partial l} \kappa = 0$$

For example, prescribe $A3$, $A4$, κ , evolve first equation for $A5$, use other two as the inhomogeneous Neumann data on $A1$ and $A2$.

The static constraints equations:

$$\partial_t^2 A_{ij} = \partial^l \partial_l A_{ij} - 2\partial_{(i} M_{j)}$$

$$\partial_t^2 \kappa = \partial^l \partial_l \kappa - \frac{3}{2} \partial^l M_l$$

Implies

$$\partial_t^2 M_i \equiv 0!!!$$

**NO NEED IN
CONSTRAINT-PRESERVING BOUNDARY
CONDITIONS!**

Linearized BSSN, densitized lapse ($\alpha = 6\varphi$, $\beta_i = 0$)

$$\partial_t \varphi = -\frac{1}{6} \kappa,$$

$$\partial_t \kappa = -6 \partial^l \partial_l \varphi,$$

$$\partial_t \tilde{\gamma}_{ij} = -2 A_{ij},$$

$$\partial_t A_{ij} = -\frac{1}{2} \partial^l \partial_l \tilde{\gamma}_{ij} + \partial_{(i} \Gamma_{j)} - 8 \partial_i \partial_j \varphi,$$

$$\partial_t \Gamma_i = -\frac{4}{3} \partial_i \kappa.$$

Constraint equations:

$$\partial^p \partial^q \tilde{\gamma}_{pq} - 8 \partial^l \partial_l \varphi = 0, \quad \text{and/or} \quad \partial^l \Gamma_l - 8 \partial^l \partial_l \varphi = 0;$$

$$\partial^j A_{ij} - \frac{2}{3} \partial_i \kappa = 0;$$

$$\Gamma_j = \partial^l \tilde{\gamma}_{lj}.$$

Energy estimate with boundaries

(C.Gundlach and J.M.Martin-Garcia)

Growth of energy

$$\begin{aligned} \epsilon = & \|\kappa\|^2 + \|A\|^2 + 36\|\partial_l\varphi\|^2 + \|\Gamma_l - 8\partial_l\varphi\|^2 \\ & + \left\| \frac{1}{2}\partial_l\tilde{\gamma}_{ji} - \delta_{l(i}(\Gamma_{j)} - 8\partial_{j)}\varphi) \right\|^2, \end{aligned}$$

is determined by three boundary terms

$$\begin{aligned} \partial_t\epsilon = & -6 \int_{\partial\Omega} \left(\frac{\partial}{\partial n}\varphi\right)\kappa - \int_{\partial\Omega} \left(\frac{\partial}{\partial n}\tilde{\gamma}_{ij}\right)A^{ij} \\ & + 2 \int_{\partial\Omega} n_{(i}(\Gamma_{j)} - 8\partial_{j)}\varphi)A^{ij}. \end{aligned}$$

Input is needed

- De Rham complex, de Rham complex, de Rham complex!
- Semigroup theory, proofs of existence in $H(\mathbf{div})$ -like spaces.
- Analysis of long term stability for nonlinear equations.
- Nonlinear energy estimates.